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1 Structural defects in 2D MoS<sub>2</sub> nanosheets and their roles in the adsorption of

2 airborne elemental mercury

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### 11 ABSTRACT

12 In this research, ab initio calculations and experimental approach were adopted to reveal the mechanism of Hg<sup>0</sup> adsorption on MoS<sub>2</sub> nanosheets that contain various types of defects. The ab 13 14 initio calculation showed that, among different structural defects, S vacancies (Vs) in the  $MoS_2$ nanosheets exhibited outstanding potential to strongly adsorb Hg<sup>0</sup>. The MoS<sub>2</sub> material was then 15 16 prepared in a controlled manner under conditions, such as temperature, concentration of precursors, etc., that were determined by adopting the new method developed in this study. 17 18 Characterisation confirmed that the MoS<sub>2</sub> material is of graphene-like layered structure with 19 abundant structural defects. The integrated dynamic and steady state (IDSS) testing demonstrated that the Vs-rich nanosheets showed excellent Hg<sup>0</sup> adsorption capability. In 20 21 addition, ab initial calculation on charge density difference, PDOS, and adsorption pathways 22 revealed that the adsorption of Hg<sup>0</sup> on the Vs-rich MoS<sub>2</sub> surface is non-activated chemisorption. Keywords: MoS<sub>2</sub> nanosheet, structural defect, Hg<sup>0</sup> adsorption, ab initio calculation 23

# 24 **1. Introduction**

The airborne elemental mercury (Hg<sup>0</sup>) is the most dominant and problematic form of mercury[1-3] because it is highly volatile and insoluble in water [4-6]. It can travel in the air over a long-distance, bio-accumulate in the ecosystem and is persistent in the environment [7]. Therefore, it has significant impacts on ecosystem and human being's health. In the past two decades, mercury emission has become a global concern and has attracted tremendous amount of attention in the development of novel materials for the
 mercury removal[8].

Technologies for the removal of elemental mercury in gas phase is highly demanded in industries such as coal-fired power plants, petrochemical processes, incineration plants[9-13]. Currently, activated carbon is commercially applied in power plants for elemental mercury removal[14]. However, the spent carbon is still a hazardous material and may compromise fly ash as a saleable by-product[15-17]. Therefore, the development of alternative non-carbon-based sorbents is necessary for mercury emission control in industry[18-20].

Recently, 2-dimensional transition-metal dichalcogenides (2D-TMDs) have attracted 39 increasing attention due to their unique structural features and have been applied in 40 various industrial processes[21]. MoS<sub>2</sub> is a 2D-TMDs material that has layered structure, 41 42 strong in-plane bonding and weak out-of-plane interactions, which shows excellent 43 electrical, physical, chemical and mechanical properties [22]. To date, MoS<sub>2</sub> has been successfully applied in various applications for reducing the emission of hazardous 44 45 pollutants[23-26]. However, not much work has been carried out on the application of  $MoS_2$  in the adsorption of elemental mercury ( $Hg^0$ ). 46

In our previous studies, a suite of Mo-based transition metal oxides and sulphides were investigated as candidate materials for Hg<sup>0</sup> removal[27, 28]. The presence of Mo species in metal oxides was found to promote Hg<sup>0</sup> removal[27], and the MoS<sub>2</sub> nanosheets exhibited outstanding Hg<sup>0</sup> removal performance [16, 29-31]. The effects of temperature, space velocity, the existence of other gas species and the recovery of Hg0 as a resource (or for permanent disposal) have been studied in detail[30, 31].However, the effect of structural defects of MoS<sub>2</sub> nanosheets on Hg<sup>0</sup> adsorption remains unclear.

Normally, structural defects provide extraordinary opportunities to tailor the intrinsic
 properties and create new functionalities. For 2D materials, structural defects have been

found to have strong influence on the electron–electron interactions. Different types of defects in MoS<sub>2</sub> nanosheets were found playing important roles in various applications[32]. However, the roles of different types of defects in MoS<sub>2</sub> nanosheets in airborne Hg<sup>0</sup> adsorption remain unexplored.

In this research, ab initio calculation was conducted to study mechanism of the adsorption of  $Hg^0$  on  $MoS_2$  nanosheets with different types of defects. A novel method was developed to determine conditions for the rapid synthesis of  $MoS_2$ -containing materials with desired functions. Characterisation of defect-rich  $MoS_2$  containing materials was then carried out, followed by the testing of the performance of this material in  $Hg^0$  adsorption.

### 66 **2. Materials and methods**

67 **2.1** Equipment

A dedicated rig was adopted for the evaluation of mercury adsorption performance of the MoS<sub>2</sub> materials. The rig consists of a mercury analyser, elemental mercury calibrator/generator, dilution probe controller, mercury sample conditioner, which was used to generate elemental mercury and to perform online analysis of elemental and oxidised mercury. A flue gas analyser was used to analyse other gas compounds in the simulated flue gas from coal-fired power stations. The details of this experimental setup are described elsewhere[27-29, 32].

#### 75 **2.2 Computational method**

ab initio calculation was performed using Density Functional Theory – Dispersion Correction (DFT-D) with exchange-correlation functional GGA-PW91-OBS[33]. The ultra soft pseudo potential (USPP)[34] was implemented. It was assumed that there were no atomic layers being constrained in the MoS<sub>2</sub> (4×4)[35]. Furthermore, the transition state (TS) confirmation was implemented following the nudged elastic band (NEB) method.

The p-band center ( $\Delta E_I$ ), following the d-band model approach [36], is adopted to study the formation of chemisorption bonds[37].

83 
$$\Delta E_I \propto \frac{V^2}{|\varepsilon_p - \varepsilon_s|}$$
 (1)

where the right side term represents the attraction between adsorbate and adsorbent, in which V is the coupling matrix element;  $\boldsymbol{\varepsilon}_s$  is the energy level of the most active band in the adsorbate and was set to be the value of Hg<sup>0</sup> s-orbital band center;  $\boldsymbol{\varepsilon}_p$  is the energy level of p orbitals, which is the active orbital of atoms and can be calculated by  $\boldsymbol{\varepsilon}_p = \frac{\int_{-\infty}^{fermi} x_p(x) dx}{\int_{-\infty}^{fermi} \rho(x) dx}$  (the x and  $\boldsymbol{\rho}(x)$  correspond to energy (eV) and DOS (electron/eV), respectively, while integral domain is from the minimum energy to fermienergy of p-orbital electrons).

91 The adsorption energy ( $E_{ads}$ ) is calculated by,

92 
$$E_{ads} = E_{Hg+MoS_2} - (E_{Hg} + E_{MoS_2})$$
 (2)

93 where,  $E_{Hg}$  is ground state energy of the free Hg<sup>0</sup> atom in a (11 Å)<sup>3</sup> supercell;  $E_{MoS_2}$  is

94 total energy of the  $(4 \times 4)$  supercell, and  $E_{Hg+MoS_2}$  is optimized total energy of the

- 95 system with a Hg<sup>0</sup> atom being adsorbed and the  $(4 \times 4)$  supercell.
- 96 Vacancy energy is calculated by,

97 
$$E_V = (E_{DFS} + E_{MA}) - E_{PFS}$$
(3)

98 where,  $E_{PFS}$  is total energy of the defect-free structure;  $E_{DFS}$  is total energy of the 99 structure having a particular type of vacancy; while  $E_{MA}$  is total energy of missing 100 atoms in the vacancies.

101 **2.3 Sample preparation** 

102 The Vs-rich MoS<sub>2</sub> nanosheets containing material was prepared using incipient wetness 103 impregnation (IWI) followed by sulphur-chemical vapour reaction (S-CVR) method[16, 104 32]. The precursor used was  $(NH_4)_6Mo_7O_{24} \bullet 4H_2O$ , while the commercial mesoporous  $\gamma$ -105  $Al_2O_3$  (V-SK Co., Ltd.) was selected as the support. The stable phases of the molybdenum 106 containing materials under different temperatures and partial pressure were predicted 107 using FactSage<sup>™</sup> 6.3 with a Phase Diagram module [38], which was then used to guide 108 the selection of appropriate preparation conditions to ensure the formation of MoS<sub>2</sub> 109 nanosheets under a controlled manner. By following the predicted conditions, the y-110 Al<sub>2</sub>O<sub>3</sub> saturated with (NH<sub>4</sub>)<sub>6</sub>Mo<sub>7</sub>O<sub>24</sub>•4H<sub>2</sub>O solution was dried in an oven at 120 °C for 24 h and calcined in air at 520 °C for 12 h. Prior to S-CVR, the sample was pre-heated to 111 500°C and kept isothermal for 1 h with nitrogen purge at a flow rate of 100 ml/min to 112 113 remove moisture and other adsorbed matters. The samples were then in situ sulphided 114 with a  $H_2S/H_2$  gas mixture (10 vol.%  $H_2$ ) at a flow rate of 20 ml/min under a desired 115 temperature for 3 h. A similar procedure for sample preparation was described in detail 116 elsewhere[16, 32].

#### 117 **2.4** Characterization of the MoS<sub>2</sub> material

X-Ray Photoelectron Spectroscope (XPS) analysis was conducted using a Kratos AXIS 118 119 Ultra DLD spectrometer with an Al  $K_{\alpha}$  radiation source at room temperature and under 120 a vacuum condition (10<sup>-7</sup> Pa)[39]. Samples were analyzed using a Raman Renishaw 121 RM2000. The pump radiation was supplied by a diode laser of 514 nm and the Raman 122 emission was focused using a 50 × objective[40]. The crystallinity of the sample was analysed using an X-Ray Diffraction (XRD, Bruker D8 A25, Germany), during which the 123  $CuK_{\alpha}$  radiation was set as 40 kV and 40 mA, and scanning was carried out from 10° to 90° 124 125 [32]. A high-resolution transmission electron microscopy (HRTEM) (JEM 2100) operated 126 at 200 kV was used to study the morphology of the samples[39].

127 In-situ Diffuse Reflectance Infrared Fourier Transform Spectroscopy (In-situ DRIFTS) was 128 applied to investigate the surface acid sites by using NH<sub>3</sub> as the alkaline probe molecule. 129 The IR spectra were recorded using a Fourier Transform Infrared Spectrometer (Bruker V70, USA) equipped with a Praying Mantis<sup>™</sup> reaction chamber (Harrick, USA). Similar 130 experimental procedures are described elsewhere[32]. Adsorption/desorption of NH<sub>3</sub> 131 132 (as the base probe molecule) on the surface of the sample was studied by using Temperature Programmed Desorption (TPD) following similar procedures as the in-situ 133 DRIFTS experiments. 134

#### 2.5 Evaluation of mercury adsorption performance 135

The samples were evaluated by using an integrated Hg<sup>0</sup>-TPSR dynamic and steady state (IDSS) 136 method. The Hg<sup>0</sup>-TPSR and IDSS were carried out using a dedicated experimental rig that 137 equipped with a mercury generator (Tekran 3310, USA) and mercury analysis system (Tekran 138 3300RS, USA) as described in our previous research[27]. 139

Based on qualitative analysis, the characteristics of Hg<sup>0</sup> capture were extracted from the profile 140

generated from the Hg<sup>0</sup>-TPSR experiment. 141

Quantitative analysis of the Hg<sup>0</sup>-TPSR experimental results was conducted to study the 142 instantaneous Hg<sup>0</sup> removal efficiency, the adsorption and desorption of Hg<sup>0</sup> on the surface of 143 the adsorbents. 144

The maximum instantaneous Hg<sup>0</sup> removal efficiency ( $\Delta X_{max}$ ) was calculated using Equation (4). 145

146 
$$\Delta X_{\max} = \frac{[Hg^0]_{in} - [Hg^0]_{out}}{[Hg^0]_{in}} \times 100$$
(4)

Area of the adsorption region (S<sub>a</sub>) and the desorption region (S<sub>d</sub>), which corresponds to the areas 147 below or above the concentration baseline, was determined by the integration of Equation (5). 148 ct2 )

149 
$$S = \int_{t1}^{t2} f(t)dt$$
 (5)

150 where,  $t_1$  is the starting time of adsorption/desorption, min;  $t_2$  is the ending time of

adsorption/desorption, min; f(t) is the change of Hg<sup>0</sup> concentration as a function of time; S is

152 the area determined using the integration method, min· $\mu$ g·m<sup>-3</sup>.

Based on the Hg<sup>0</sup>-TPSR dynamic transient analysis, a steady-state evaluation of Hg<sup>0</sup> adsorption performance was further conducted. The samples were evaluated under the same temperature (50°C) and Hg<sup>0</sup> concentration (30  $\mu$ g/m<sup>3</sup>) for at least 180 min after the reactions had reached steady state.

157 **3. Results and discussion** 

#### 158 **3.1 Computational study**

### 159 **3.1.1** Geometric structure of defective MoS<sub>2</sub> nanosheets

160 The MoS<sub>2</sub> nanosheets normally contain different structural defects in its crystal lattices, 161 which are the potential sites for airborne  $Hg^0$  capture [41].The MoS<sub>2</sub> nanosheets with 162 defect-free surface as well as those with different types of defects, such as vacancies 163 with the absence of one S atom (V<sub>S</sub>), the absence of two S atoms (V<sub>S2</sub>) and the absence 164 of one Mo atom (V<sub>Mo</sub>), were investigated in detail in this study.

165 It is found that the atoms around the S vacancy exhibit different extents of inward 166 relaxation depending on the number of S atoms that are absent. This is because the 167 radius of an S<sup>2-</sup> ion (0.184nm) is much larger than that of a Mo<sup>4+</sup> ion (0.065nm) [42, 43] 168 and this type of inward relaxation is commonly seen in defective surface structures with 169 large size atoms being removed from a defect-free surface.

For a Mo vacancy, its neighbouring atoms show an outward relaxation due to the weakened electrostatic attraction force between the NA S atom and the Mo vacancy. Additionally, the outward relaxation of the NA S atom is greater than that of the NA Mo atoms. That means the relative distance between upper S plane and the second NA Mo atoms becomes greater as compared with no relaxation occurred. Based on the calculation, after relaxation, the distance between the NA S atoms and the Mo vacancy increased by 2.61%, while the distance between the NA Mo atoms and the Mo vacancy was shorten by 0.63% (as shown in Table 1). It can be concluded that absence of Mo atom(s) leads to the formation of Mo vacancies that are negative charged, which have the potential to facilitate the adsorption of metals in oxidized state. Therefore, the oxidation of Hg species (for example, HgCl, HgO) could have affinity with such kind of positions.

182 However, driven by the weakened electrostatic repulsion between the nearest adjacent 183 (NA) S atoms and the S vacancy, the extent of the inward relaxation of the NA S atom is 184 greater than that of the NA Mo atoms. Therefore, the distances between V<sub>s</sub> and the NA 185 Mo atoms and between  $V_s$  and the NA S atoms are shortened by 0.04% and 2.70%, 186 respectively. Therefore, it can be concluded that the absence of S atoms leads to the 187 formation of S vacancies that are positive charged (electron deficient), which has the 188 potential to facilitate the capture of metal atoms (electron donors). Moreover, it is found that the inward relaxation of the  $V_{S2}$  is greater than that of the  $V_S$  as the  $V_{S2}$ -Mo<sub>NA</sub> and 189 190 the V<sub>S2</sub>-S<sub>NA</sub> distances are shortened by 5.47% and 3.53% respectively (Table 1). This 191 means that after relaxation, distance between the Mo atoms and the S vacancy become 192 shorter.

Therefore, the missing of S atoms leads to the formation of S vacancies that are positive charged (electron deficient), which would facilitate the capture of metal atoms (electron donors). Moreover, some researches have reported the adsorption of metals, such as Mn atom, Fe atom and Au cluster, on the Vs type of defects in MoS<sub>2</sub> nanosheet [44, 45], which shows the potential of Vs position in capturing metal atoms. Most of the trace Fe and Mn species were in dust form of flue gas at coal fired power plant that can be removed by existing air pollution control devices (APCDs) such as electrostatic

- 200 precipitator (ESP), fabric filters (FF). Therefore, most of the effective surface of the MoS<sub>2</sub>
- 201 with Vs defect is active for capture airborne Hg<sup>0</sup>.

### 202 **3.1.2 ab initio mapping of MoS<sub>2</sub> surface**

Theoretically, a more negative adsorption energy ( $E_{ads}$ ) indicates a more favourable interaction between adsorbent and the adsorbate [46]. All potential adsorption sites on the defect-free and defective MoS<sub>2</sub> nanosheets were evaluated and numbered in the order of magnitude of the adsorption energy as shown in **Figure 1**.

Based on ab initial calculations, it is clear that on the defect-free  $MoS_2$  nanosheets, the strongest adsorption site is Position 18, which is the position on top of a Mo atom. This suggests that the adsorption of  $Hg^0$  on the  $MoS_2$  nanosheets is affected by S and Mo atoms in a synergistic manner.

211 For the defective structure with one absent Mo atom  $(V_{Mo})$ , the site with the second highest absolute adsorption energy (Position 21) is also found to be on top of a Mo atom. 212 213 However, for this V<sub>Mo</sub> structure, the strongest adsorption site is on top of the hollow space of the sulphur plane (Position 2). This is mainly due to the repulsive force against 214  $Hg^0$  atom from the negatively charged  $V_{Mo}$ . With the absence of Mo atom, the adsorption 215 energy of  $V_{Mo}$  is significantly lower than that of the other positions on the defective 216 surfaces and the defect-free surface. This further indicates that Mo atoms affect the 217 adsorption of Hg<sup>0</sup> on MoS<sub>2</sub> nanosheets. 218

Based on ab initio calculations, it is clear that  $Hg^0$  adsorption occurred on the MoS<sub>2</sub> nanosheets with V<sub>S2</sub> and V<sub>S</sub> showed similar behaviours. For both cases, the adsorption site with the second largest absolute adsorption energy, Position 2 for V<sub>S</sub> and Position 11 for V<sub>S2</sub>, are on top of Mo atoms, with an adsorption energy of approximately 1.12 eV. Furthermore, for the adsorption of mercury atoms on defective MoS<sub>2</sub> nanosheets, i.e., V<sub>S2</sub> and V<sub>S</sub>, the most stable adsorption sites are the ones right above the original position of the absent S atoms, i.e., on top of the S vacancies. It is found that for the structure with only one absent sulphur atom, the absolute adsorption energy is slightly greater than that of the structure with two absent sulphur atoms, which indicates stronger attraction to  $Hg^0$  on the V<sub>s</sub> than on the V<sub>s2</sub>. These results suggest that the S vacancies have major influence on the adsorption of  $Hg^0$ , while the absence of an additional S atom in the V<sub>s2</sub> has less influence on the adsorption process.

In order to quantify the interactions between Hg<sup>0</sup> and the different types of vacancies on
a MoS<sub>2</sub> surface, p-band centre analysis was carried out following a similar method as the
d-band centre analysis for transition metals [37]. Compared with PDOS or DOS of each
element in MoS<sub>2</sub> as shown in

Figure 2, it can be found that both the s and p orbitals of S atoms and s and p orbitals of 235 Mo atoms have some noticeable influence on the adsorption of Hg<sup>0</sup>. Moreover, p-band 236 237 centre analysis was also carried out in this study to show the bonding strength of mercury adsorption on different types of defects in MoS<sub>2</sub> surface. The variation in binding energy 238 (Equation (1)) is determined by the  $|\varepsilon_p - \varepsilon_s|$ , which is plotted in **Figure 3** (a). The closer 239 of p band center to the Fermi level indicates the greater interaction between s orbital of 240 the Hg<sup>0</sup> atom with p orbitals of the defective MoS<sub>2</sub> surface. Based on p-band center 241 analysis, it is found that adsorption energy is in the order of V<sub>S</sub> > V<sub>S2</sub> > V<sub>Mo</sub>. Therefore, it 242 can be concluded that V<sub>s</sub> is the strongest sites for Hg<sup>0</sup> adsorption. 243

Generally, defects with different formation energy can affect the physical and chemical properties of the MoS<sub>2</sub> monolayers, which could subsequently influence adsorption[47]. To further reveal the adsorption of Hg<sup>0</sup> on different defects in the MoS<sub>2</sub> surfaces, relationship between formation energy ( $E_V$ ) of the defects and adsorption energy was studied using Equations (2) and (3), the results of which are shown in **Figure 3** (b). The smaller positive E<sub>V</sub> value suggests that the formation of the vacancy is endothermic and easier to form. As shown in Figure 3 (b), the V<sub>s</sub> had the lowest formation energy. This is consistent with findings of other researchers that the V<sub>s</sub> is frequently observed in MoS<sub>2</sub> surfaces [48]. The formation energy of V<sub>s2</sub> was higher than that of V<sub>s</sub>, which is consistent with reported finding that randomly distributed V<sub>s</sub> is more frequently observed than V<sub>s2</sub> [48]. The formation energy of V<sub>Mo</sub> is even higher than that of V<sub>s2</sub>. These results further indicate that the Vs is the most stable configuration for Hg<sup>0</sup> adsorption and the Vs is the prevailing defect in MoS<sub>2</sub> nanosheets.

#### 257 3.2 Experimental Study

#### 258 **3.2.1 Preparation of defective MoS<sub>2</sub> nanosheets**

In this study, a novel approach based on FactSage<sup>™</sup> (version 6.3) with phase diagram 259 module was developed to determine conditions for sample preparation, such as 260 261 concentration of precursors and temperature for preparation, which is to avoid problems 262 associated with traditional trial-and-error approach in sample preparation. Figure 4 (a) 263 illustrates preparation conditions for the formation of MoS<sub>2</sub> sample via sulphur-chemical vapour reaction (S-CVR) process. This triangle phase diagram indicated that the calcined 264 265 material (mainly consists of MoO<sub>3</sub>) reacts with H<sub>2</sub>S and H<sub>2</sub> gas mixture with different mole fractions at 400 °C and results in the formation of various stable phases, such as MoS<sub>2</sub>, 266 267  $MoO_2$ ,  $Mo_4O_{11}$ ,  $Mo_8O_{23}$  and  $Mo_9O_{26}$ . It is clear that to prepare  $MoS_2$  as the main component of the adsorbent, the calcined adsorbent should be sulphided at 400 °C in 268 269 the atmosphere of H<sub>2</sub>S (10 vol.%) /H<sub>2</sub> (90 vol.%). It can be concluded that the graphene-270 like MoS<sub>2</sub>-containing adsorbent (with a surface area around 204  $m^2/g$ ) was prepared on the mesoporous  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (with 211 m<sup>2</sup>/g surface area) in a well-controlled manner. This 271 272 novel method developed in this study can also be used in the controlled synthesis of 273 other materials for environmental applications.

#### 274 **3.2.2** Characterization of Defective MoS<sub>2</sub> Nanosheets

275 The samples prepared were then characterized by using a number of techniques. The 276 well-dispersed chemical species and the chemical state on the surface of samples were 277 further investigated by using XPS analysis[49, 50]. The stable phases of the Mo 3d XPS spectra are as shown in Figure 4(b). According to the XPS database (version 4.1) of NIST 278 Chemistry WebBook [51], the peaks with binding energy (BE) of approximately 232.8 eV, 279 280 228.8 eV, 162.6 eV are assigned to  $Mo^{4+}$  3d3/2,  $Mo^{4+}$  3d5/2 and S 2p, respectively. 281 Therefore, it is proved that MoS<sub>2</sub> was successfully prepared in this study. This also confirmed that it is feasible to use FactSage<sup>™</sup> to guide the preparation of samples with 282 desired phase of the active components. 283

The S-Mo-S layered structure was then characterized by Raman spectroscopy. As shown 284 in **Figure 4**(c), the Raman spectra show two strong peaks at the in-plane  $E_{2g}^{1}$  and the out-285 286 of-plane A<sub>1g</sub> vibration, which are the typical peaks of MoS<sub>2</sub> nanosheets. Normally the 287 peak positions of E<sup>1</sup><sub>2g</sub> and A<sub>1g</sub> bands are the strong indicator of the number of layers [52]. Figure 4 (c) shows the peak positions of  $E_{2g}^{1}$  and  $A_{1g}$  bands, which are 378.72cm<sup>-1</sup> and 288 404.11cm<sup>-1</sup>, respectively. The peak position of A<sub>1g</sub> bands (out-of-plane vibration, see the 289 290 schematic diagram inside the Figure 4(c)) is consistent with that of 1-3 layers. However, the frequency of E<sup>1</sup><sub>2g</sub> bands (i.e., in-plane vibration) decreased when compared with 291 292 what was reported for pristine  $MoS_2$  (384 cm<sup>-1</sup>)[52]. The significant shift of  $E^{1}_{2g}$  mode and 293 almost unchanged A<sub>1g</sub> mode correspond to sulphur vacancies in MoS<sub>2</sub> [47].

294 Moreover, XRD analysis was conducted to show structure of the MoS<sub>2</sub> on  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> 295 support. Only three major peaks at around 37, 47 and 67 degrees can be identified as 296 shown in **Figure 4**(d) (ii), which are the characteristic peaks of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> crystal (as shown 297 in **Figure 4**(d)(i)). There are no visible diffraction peaks of Mo species, which is an 298 indication that MoS<sub>2</sub> is well-dispersed on the support and forms monolayers 299 (nanosheets) [53].

Morphology of the adsorbent prepared was further characterized using a HRTEM as 300 301 shown in **Figure 4**(e). It can be seen that the regular interplanar spacing is 0.27 nm. This 302 is the characteristic pattern of (001) planes of the hexagonal MoS<sub>2</sub> crystal structure. Based on this observation, the schematic structure of the 2D-TMD MoS<sub>2</sub> nanosheets was 303 304 proposed and illustrated inside the **Figure 4**(e). The HRTEM result is consistent with the other researchers' recent findings that the un-supported MoS<sub>2</sub> sample is defect-rich 305 ultra-thin MoS<sub>2</sub> nanosheets [48]. In addition, there are many dislocations and distortions 306 307 that can be observed in Figure 4(e), which also suggest that this type of MoS<sub>2</sub> nanosheets contains defects [48]. 308

309 Moreover, NH<sub>3</sub> was selected as the probe molecule for the in-situ DRIFTS to characterize acidity of active centres (Figure 4(f)), which is to show the surface property and therefore 310 311 to understand properties of the defect-rich MoS<sub>2</sub> nanosheet [52]. The peaks at around 312 1230, 1450, 1630, 1670, 2800, 3000 and 3200 cm<sup>-1</sup> remained after NH<sub>3</sub> adsorption and N<sub>2</sub> purging. Normally, peaks at 1230 cm<sup>-1</sup> and 1630 cm<sup>-1</sup> are the characteristic pattern of 313 314 the absorbed NH<sub>3</sub> molecules on Lewis acid sites, whilst peaks at 1450 cm<sup>-1</sup> and 1670 cm<sup>-1</sup> 315 <sup>1</sup> are the coordinated NH<sub>4</sub><sup>+</sup> species on Brønsted acid sites [54]. The peaks at 2800, 3000 and 3200 cm<sup>-1</sup> in the stretching region are associated with the NH<sub>4</sub><sup>+</sup> bands. It is obvious 316 317 that the intensity of Brønsted acid sites of the MoS<sub>2</sub> material was higher than that of Lewis acid sites. 318

Surface acidity was further investigated by using NH<sub>3</sub>-TPD. The NH<sub>3</sub> desorption peaks at low temperatures are attributed to the Brønsted acid sites, whereas those peaks at high temperatures are assigned to Lewis acid sites [27]. As shown in **Figure 4**(f), the strong desorption peaks are observed at temperature below 250 °C, which further confirms that surface of the MoS<sub>2</sub> nanosheets is dominated by Brønsted acid sites. The amount of NH<sub>3</sub> adsorbed at Brønsted acid sites can be associated with the existence of coordinative unsaturated sites (CUS) on surface of the defect-rich MoS<sub>2</sub> nanosheets [55]. The 13 326 formation of Vs on the MoS<sub>2</sub> nanosheets might result in high Hg<sup>0</sup> adsorption efficiency

327 as previously predicted by computational results.

# 328 **3.2.3** Evaluation of Defective MoS<sub>2</sub> Nanosheets for Hg<sup>0</sup> Adsorption

Mercury adsorption capability of the defective MoS<sub>2</sub> nanosheets was then evaluated by using the IDSS experimental method. **Figure 5**(a) shows that the adsorbent demonstrated excellent performance in the adsorption of Hg<sup>0</sup> at temperatures below 150°C. The steady-state analysis (as shown in **Figure 5**(b)) further demonstrates that the Vs-rich MoS<sub>2</sub> nanosheets show excellent performance in Hg<sup>0</sup> adsorption with almost 100% removal efficiency.

It is found that the adsorbed Hg<sup>0</sup> could then be desorbed when temperature was raised 335 to above 200°C with the calculated ratio of desorption over adsorption exceeding 99.4%. 336 337 This suggests that the MoS<sub>2</sub> nanosheet containing materials can be regenerated easily at 338 high temperature. This property makes MoS<sub>2</sub> nanosheet containing material a better adsorbent both technically and economically, as compared with activated carbon 339 injection technology. In addition, this Vs-rich MoS2 nanosheets demonstrates the 340 potential to recover the captured Hg<sup>0</sup> efficiently as a resource without having any 341 impacts on the saleability of fly ash as a product. 342

## 343 **3.3 Hg<sup>0</sup> adsorption mechanism**

Mechanism of the adsorption of Hg<sup>0</sup> on the Vs-rich MoS<sub>2</sub> surface was further studied by the charge density difference analysis, PDOS analysis, and adsorption pathways and energy profiles predictions.

Plots of the charge-density difference for  $Hg^0$  adsorption on the Vs, Vs<sub>2</sub> and V<sub>Mo</sub> are shown in **Figure 6**. The blue area covered the entire S vacancy, which means charge density increased in this area.

350 Atomic charges for Vs in MoS<sub>2</sub> monolayer and corresponding Hg<sup>0</sup>-MoS<sub>2</sub> nanosheets were

351 calculated by using Mulliken method. Table 2 shows the electron transfer through Hg<sup>0</sup>-

MoS<sub>2</sub> (001) interactions with Vs, Vs<sub>2</sub> and V<sub>Mo</sub>. The charge transfer was based on the difference between the original state and the adsorption state. A positive value means the atom obtains electrons after adsorption, while a negative value indicates the atom loses electrons after adsorption.

Results of Mulliken charge analysis (as illustrated in Table 2) showed that the electrons 356 transferred a charge of 0.02eV from the Hg atom (6s) to S atom (3p) through the sulphur 357 vacancy (Vs), and the Mo atom (4p) also transfers a charge of 0.01eV to the S surface, 358 which enhanced the electronegativity of the vacancy (shown in blue area). There are 359 charge-charge interactions between Hg<sup>0</sup> and the Vs, which indicates that Hg<sup>0</sup> atoms are 360 chemically adsorbed on the surface. The order of the magnitude of adsorption energy at 361 Vs (as shown in Figure 1) also corresponds to chemical interactions between the surface 362 363 and the mercury atoms adsorbed [56]. The distance from the mercury atom to the Vs is short as shown in Figure 7. Since the distance from the Hg<sup>0</sup> atom to the NA S atoms is 364 365 shorter than the sum of the covalent radii by 40%[57], it is confirmed again that the Hg atom is chemically adsorbed on the Vs. 366

To further understand the adsorption of  $Hg^0$  atoms on  $MoS_2$  (Vs) nanosheets, the density of state (DOS) of the surface atoms was studied in depth. The partial density of states (PDOS) of Hg, Mo and S atoms were calculated and are shown in

Figure 2. The PDOS peaks of d, s and p orbitals of an isolated Hg<sup>0</sup> atom are close to -3.1, 0 and 5.7eV, respectively. After adsorption, all the PDOS peaks of the Hg atom shifted left with the state of s and p orbitals significantly decreased in energy level. This suggested that there are strong interactions between Hg<sup>0</sup> and MoS<sub>2</sub> (Vs-rich) nanosheets. As shown in

Figure 2 (a), no noticeable shift in d states of the Mo atom was observed, which suggests that 4d orbitals of Mo atoms does not have noticeable influence on the adsorption of Hg<sup>0</sup>. Among these notable features of the PDOS changes, it is worth mentioning that the

peak at -4 eV of PDOS(d) of Hg atoms overlaps with the major peak in the PDOS(s) of Mo 378 379 atoms and PDOS(s) of sulphur atoms. This indicates the repulsive force between Hg 380 atoms and MoS<sub>2</sub> nanosheets may be from the filled s orbitals of S, d orbitals of Hg and unfilled s orbitals of Mo. Compared with the PDOS of p states of Mo and S atoms before 381 382 and after adsorption, both strongly overlapped with the s and p states of Hg atoms at approximately -0.9 eV and 2.3 eV (the value of DOS, 0.1/ eV), respectively. All of these 383 suggest that Hg<sup>0</sup> atoms interact strongly with the MoS<sub>2</sub> surface, and the p states of both 384 385 Mo and S atoms play important roles in the adsorption of mercury. This is also the reason 386 why the p-band centre analysis was developed for potential adsorption sites evaluation. Figure 8 illustrates the adsorption pathways and energy profiles of a Hg<sup>0</sup> atom adsorbed 387 on the Vs in MoS<sub>2</sub> (001) surface. The result represents the minimum energy path for the 388 389 diffusion of an Hg<sup>0</sup> atom between Position 9 and Position 1 as shown in Figure 1.

Based on the fully optimized adsorption geometry of the Vs of MoS<sub>2</sub> (001) nanosheet 390 391 (monolayer) (as shown in Figure 1(a)), nine stable adsorption configurations were 392 identified. As discussed, Position 1 corresponds to a Hg atom adsorbed on top of the S 393 vacancy, which is the most stable adsorption site. Position 9 represents the adsorption of a Hg atom on top of an S atom, which is the second NA S atom to the vacancy with the 394 395 minimum adsorption energy. The minimum energy paths for the diffusion of an Hg atom between different adsorption configurations using NEB calculations were further 396 397 investigated and are shown in Figure 8.

The NA Mo atoms around the S vacancy make the Hg adsorption configuration unstable, which is caused by the interactions between the Hg cations and the Mo cations. Furthermore, the positive charge (electron deficient) of Vs attracts the Hg atom to the centre of the vacancy. As discussed previously, the inward relaxation makes the surface of MoS<sub>2</sub> become a sunken cave. As shown in Figure 8, the pathway between Position 9 and Position 1, from high energy level to low energy level, has no energy barrier for the

adsorption to occur. Based on the adsorption pathways and energy profiles predictions, the adsorption could be classified as non-activated chemisorption, which is introduced by surface defects (V<sub>S</sub>), which exhibit high adsorption probability[58]. The IDSS experimental observation indicated that the Vs containing MoS<sub>2</sub> nanosheets are able to adsorb Hg<sup>0</sup> at lower temperatures, which is consistent with ab initio calculation computational investigation. It can therefore be concluded that the adsorption of Hg<sup>0</sup> atom around the Vs is non-activated chemisorption.

#### 411 **4. Conclusions**

In this study, the ab initio calculation showed that S vacancies (Vs) in the MoS<sub>2</sub> 412 nanosheets exhibited outstanding potential in Hg<sup>0</sup> adsorption among different structural 413 defects. The MoS<sub>2</sub> material was successfully prepared in a controlled manner under 414 415 conditions that were determined via phase diagrams that are generated using Factsage. 416 Results of characterisation work confirmed that the MoS<sub>2</sub> material consists of nanosheets and is of graphene-like layered structure with abundant defects. Moreover, 417 the existence of coordinative unsaturated sites (CUS) on the surface of the defect-rich 418 419 MoS2 nanosheets was confirmed by the NH<sub>3</sub>-TPD and in-situ DRIFTS by using NH<sub>3</sub> as a probe molecule. Moreover, the experimental results confirmed that this Vs-rich 420 nanosheets showed excellent Hg<sup>0</sup> adsorption. The adsorption of Hg<sup>0</sup> on Vs was found to 421 be non-activated chemisorption. 422

423

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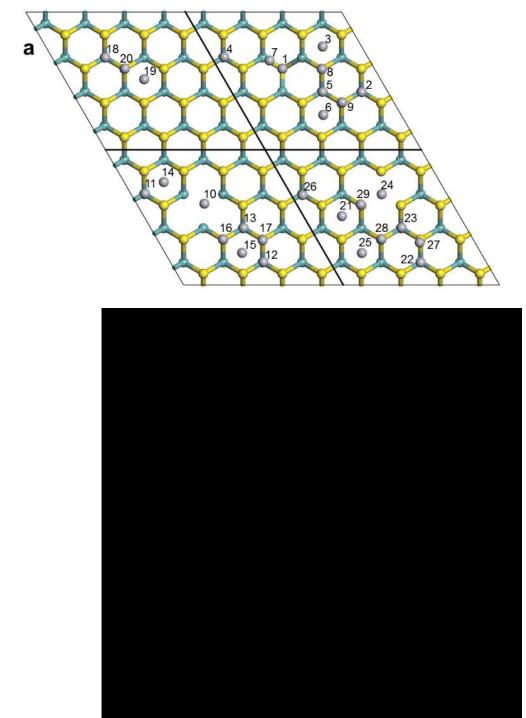
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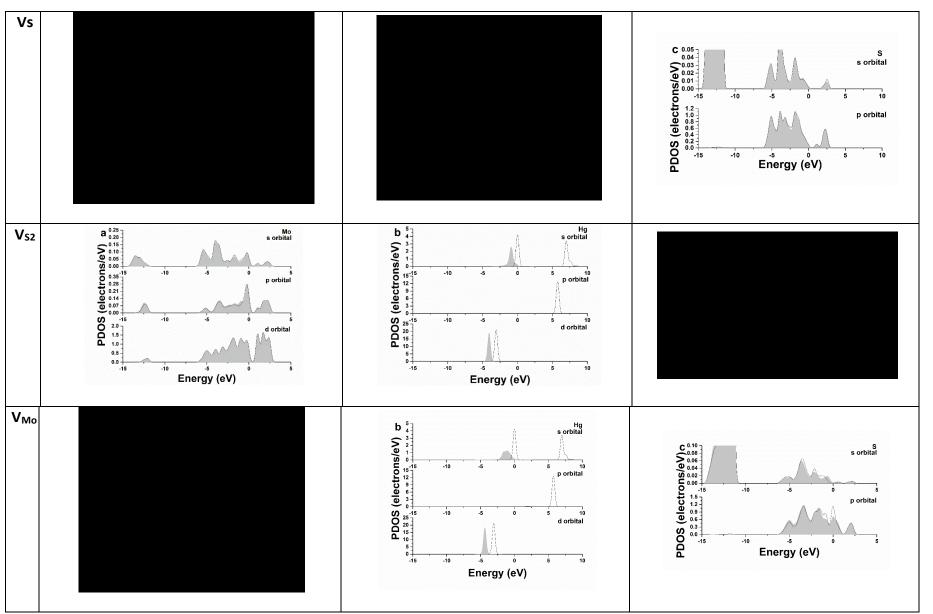
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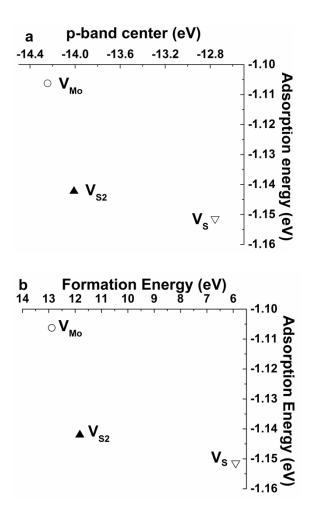
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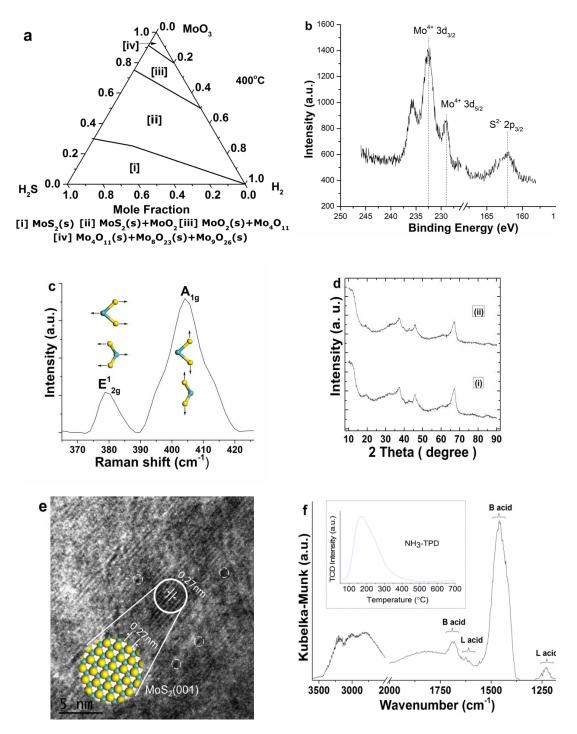
- **Figure 1** Potential Hg<sup>0</sup> adsorption sites in the MoS<sub>2</sub> surface and their corresponding
- adsorption energy. (a) Hg<sup>0</sup> adsorption sites and (b) Hg<sup>0</sup> adsorption energy



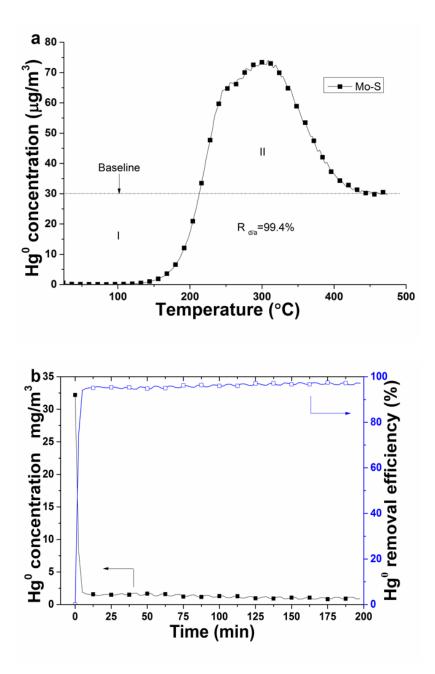
**Figure 2** DOS analysis of a defective MoS<sub>2</sub> monolayer. (a) DOS of s orbital after adsorption, (b) DOS of p orbitals after adsorption, and, (c) DOS of d orbitals after adsorption. (Black line represents V<sub>S2</sub>, the grey shadow represents V<sub>S</sub>, red line represents V<sub>Mo</sub>.



**Figure 3** Different defects for Hg<sup>0</sup> capture. (a) p-band center analysis, and, (b) formation energy of defect for the most stable Hg<sup>0</sup> adsorption sites



**Figure 4** Preparation and characterization of  $MoS_2$  nanosheets. (a) Phase diagram. (b) XPS spectra of Mo 3d and S 2p. (c) Raman spectra. (d) XRD analysis of  $\gamma$ -Al<sub>2</sub>O<sub>3</sub> (i) and the MoS<sub>2</sub> containing adsorbent(ii). (e) HRTEM picture and schematic structure (enlarged image). (f) In-situ DRIFTS study and the NH<sub>3</sub>-TPD profile (inserted figure).



**Figure 5**  $Hg^0$  adsorption performance of Vs  $MoS_2$  (001) containing adsorbent. (a) Dynamic transient and (b) steady-state experimental data

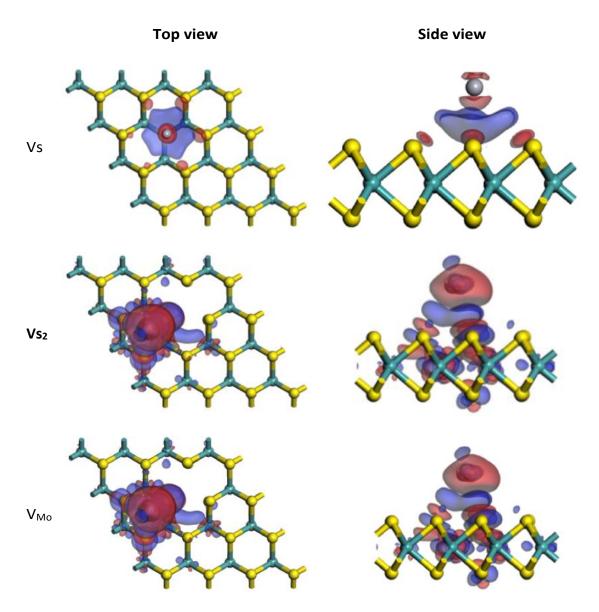


Figure 6 Charge density difference analysis of an  $Hg^0$  atom adsorbed on different types of defects in  $MoS_2$  (001) of the most stable configuration (iso-surface: 7.5×10-4 e/Å<sup>3</sup>).

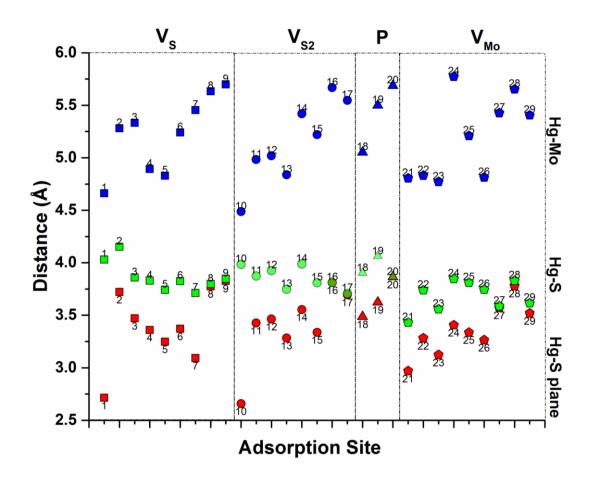


Figure 7 Hg<sup>0</sup> adsorption sites and the distances for different scenarios: Blue: distanceto its adjacent Mo atom; Green: distance to its adjacent S atom; Red: the distance to the upper layer of S atoms.

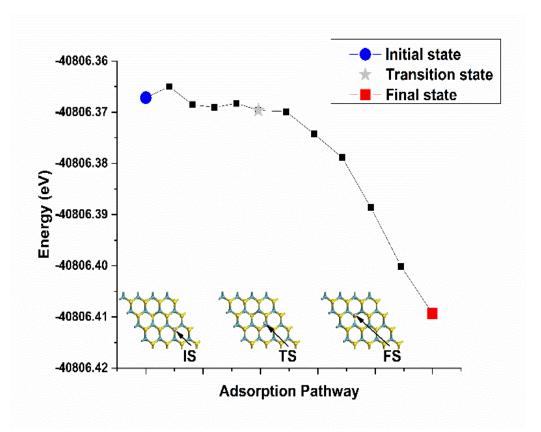


Figure 8 Adsorption pathways and energy profiles of the adsorption of an  $Hg^0$  atom on the V<sub>s</sub> MoS<sub>2</sub> (001) Surface.

# Table 1 Distances from the vacancy to the NA atoms for MoS2 monolayer before

Vacancy	Distance from vacancy to NA atoms(Å)							
	Along the S-Mo bond	δ%	Along the S-Mo bond	δ%				
	direction		direction					
Vs	2.41(V- NA Mo)	-0.04%	3.097 (V- NA S)	-2.70%				
Vs2	2.279(V- NA Mo)	-5.47%	3.071 (V- NA S)	-3.52%				
Vmo	2.474(V- NA S)	2.61%	3.163 (V- NA Mo)	-0.63%				
Perfect	2.411		3.183					

# and after geometry optimization.

\* d(Mo-Mo)=d(S-S)= 3.183 Å

Atom indexes	Vs		Atom	V <sub>s2</sub>			Atom	V <sub>Mo</sub>			
	Before	After	Electron Transfer	indexes	Before	After	Electron Transfer	indexes	Before	After	Electron Transfer
S(9)	-0.02	-0.01	-0.01	S(1)	-0.03	-0.02	-0.01	S(1)	-0.04	-0.03	-0.01
S(11)	-0.02	-0.01	-0.01	S(12)	-0.03	-0.02	-0.01	S(7)	-0.02	-0.01	-0.01
S(17)	-0.02	-0.01	-0.01	S(18)	-0.03	-0.02	-0.01	S(17)	-0.02	-0.01	-0.01
S(21)	-0.04	-0.03	-0.01	S(28)	-0.03	-0.02	-0.01	S(23)	-0.02	-0.03	0.01
S(26)	-0.02	-0.01	-0.01	S(29)	-0.03	-0.02	-0.01	S(24)	-0.02	-0.01	-0.01
S(28)	-0.02	-0.01	-0.01	Mo(2)	0.11	0.1	0.01	S(26)	-0.04	-0.03	-0.01
Mo(3)	0.05	0.06	-0.01	Mo(5)	0.01	-0.01	0.02	S(29)	-0.1	-0.09	-0.01
Mo(7)	0.04	0.02	0.02	Mo(10)	0.06	0.03	0.03	Mo(6)	0.1	0.07	0.03
Mo(13)	0.07	0.08	-0.01	Mo(14)	0.11	0.1	0.01	Mo(12)	0.07	0.08	-0.01
Mo(15)	0.04	0.02	0.02	Mo(15)	0.01	-0.01	0.02	Mo(14)	0.1	0.09	0.01
Hg	0	0.02	-0.02	Hg	0	0.02	-0.02	Hg	0	0.02	-0.02

# Table 2 Charge transfer of atoms for different defects determined by Mulliken Charge Analysis